

## Navier-Stokes Equations in Spherical Coordinates

In spherical coordinates,  $(r, \theta, \phi)$ , the Navier-Stokes equations of motion for an incompressible fluid with uniform viscosity are:

$$\rho \left[ \frac{Du_r}{Dt} - \frac{u_\theta^2 + u_\phi^2}{r} \right] = -\frac{\partial p}{\partial r} + f_r + \mu \left[ \nabla^2 u_r - \frac{2u_r}{r^2} - \frac{2}{r^2} \frac{\partial u_\theta}{\partial \theta} - \frac{2u_\theta \cot \theta}{r^2} - \frac{2}{r^2 \sin \theta} \frac{\partial u_\phi}{\partial \phi} \right] \quad (\text{Bhh1})$$

$$\rho \left[ \frac{Du_\theta}{Dt} + \frac{u_\theta u_r}{r} - \frac{u_\phi^2 \cot \theta}{r} \right] = -\frac{1}{r} \frac{\partial p}{\partial \theta} + f_\theta + \mu \left[ \nabla^2 u_\theta + \frac{2}{r^2} \frac{\partial u_r}{\partial \theta} - \frac{u_\theta}{r^2 \sin \theta \sin \theta} - \frac{2 \cot \theta}{r^2 \sin \theta} \frac{\partial u_\phi}{\partial \phi} \right] \quad (\text{Bhh2})$$

$$\rho \left[ \frac{Du_\phi}{Dt} + \frac{u_\phi u_r}{r} + \frac{u_\theta u_\phi \cot \theta}{r} \right] = -\frac{1}{r \sin \theta} \frac{\partial p}{\partial \phi} + f_\phi + \mu \left[ \nabla^2 u_\phi + \frac{2}{r^2 \sin \theta} \frac{\partial u_r}{\partial \phi} - \frac{u_\phi}{r^2 \sin \theta \sin \theta} + \frac{2 \cot \theta}{r^2 \sin \theta} \frac{\partial u_\theta}{\partial \phi} \right] \quad (\text{Bhh3})$$

where  $u_r, u_\theta, u_\phi$  are the velocities in the  $r, \theta, \phi$  directions,  $p$  is the pressure,  $\rho$  is the fluid density and  $f_r, f_\theta, f_\phi$  are the body force components. The Lagrangian or material derivative is

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + u_r \frac{\partial}{\partial r} + \frac{u_\theta}{r} \frac{\partial}{\partial \theta} + \frac{u_\phi}{r \sin \theta} \frac{\partial}{\partial \phi} \quad (\text{Bhh4})$$

and the Laplacian operator is

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin \theta \sin \theta} \frac{\partial^2}{\partial \phi^2} \quad (\text{Bhh5})$$

Moreover, for an incompressible fluid the equation of continuity in spherical coordinates is

$$\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 u_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (u_\theta \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial u_\phi}{\partial \phi} = 0 \quad (\text{Bhh6})$$

and for an incompressible, Newtonian fluid

$$\sigma_{rr} = -p + 2\mu \frac{\partial u_r}{\partial r} ; \quad \sigma_{\theta\theta} = -p + 2\mu \left( \frac{1}{r} \frac{\partial u_\theta}{\partial \theta} + \frac{u_r}{r} \right) \quad (\text{Bhh7})$$

$$\sigma_{\phi\phi} = -p + 2\mu \left( \frac{1}{r \sin \theta} \frac{\partial u_\phi}{\partial \phi} + \frac{u_r}{r} + \frac{u_\theta \cot \theta}{r} \right) \quad (\text{Bhh8})$$

$$\sigma_{r\theta} = \mu \left( r \frac{\partial}{\partial r} \left( \frac{u_\theta}{r} \right) + \frac{1}{r} \frac{\partial u_r}{\partial \theta} \right) ; \quad \sigma_{r\phi} = \mu \left( \frac{1}{r \sin \theta} \frac{\partial u_r}{\partial \phi} + r \frac{\partial}{\partial r} \left( \frac{u_\phi}{r} \right) \right) \quad (\text{Bhh9})$$

$$\sigma_{\theta\phi} = \mu \left( \frac{1}{r \sin \theta} \frac{\partial u_\theta}{\partial \phi} + \frac{\sin \theta}{r} \frac{\partial}{\partial \theta} \left( \frac{u_\phi}{\sin \theta} \right) \right) \quad (\text{Bhh10})$$